

Experimental data substituted in the latter equation show that, over the range 50–225 kg/cm<sup>2</sup>,  $\Delta\beta$  has a constant value of  $\sim 2.5 \times 10^{-6}$  cm<sup>2</sup>/kg, which is at most 2.5 percent of  $\beta_f$ . At 3555 kg/cm<sup>2</sup>  $\Delta\beta$  is calculated to be  $2.8 \times 10^{-6}$  cm<sup>2</sup>/kg or 5 percent of  $\beta_f$ . The compressibility coefficient of solid He<sup>3</sup> is therefore very similar to that of the fluid along the full range of the melting curve investigated. For Na and K, the data of Bridgman (34) lead to values of 38 and 29 percent, respectively, for  $\Delta\beta/\beta_f$  at  $P_m = 1$  kg/cm<sup>2</sup>.

### C. THERMAL PROPERTIES OF MELTING

At the lower end of the  $P_m$  range for He<sup>3</sup>, the  $\Delta S_m$  results were combined with the entropy of saturated liquid  $S_{\text{sat}}$ , measured by Roberts and Sydoriak (35), and the entropy of compression  $\Delta S_{\text{comp}}$  to give the entropy of solid. The values of  $\Delta S_{\text{comp}}$  can be obtained through the formula

$$\Delta S_{\text{comp}} = - \int_{P_{\text{sat}}}^{P_m} \left( \frac{\partial V}{\partial T} \right)_P dP.$$

For the computation, the present measurements were used from 5 kg cm<sup>2</sup> to  $P_m$ , and those of Sherman and Edeskuty (29), from  $P_{\text{sat}}$  to 5 kg cm<sup>2</sup>. The results over 1.2° to 2.0°K showed the entropy of solid at the melting curve (or  $S_a$ ) to rise only from 1.34 to 1.43 cal/deg mol. Subtraction of the entropy change of compression and of transition in solid gave approximate  $S_b$  values of 1.32 to 1.34. The entropy associated with a nuclear spin system in completely random orientation is  $S = R \ln 2 = 1.38$ . It would appear that for solid He<sup>3</sup> this is the major source of entropy.

The values of  $\Delta S_m$  listed in Tables I and II were derived from the Clapeyron equation using experimental  $\Delta V_m$  data and slopes computed from analytical expressions for the melting curves. For both He isotopes  $\Delta S_m$  increases with  $P_m$  over the experimental range covered, although the increase becomes progressively smaller at higher melting pressures. This behavior is contrary to that of N<sub>2</sub> (15), which showed a decrease of  $\Delta S_m$  with increasing  $P_m$ . Ebert (36), using melting properties for almost all materials studied to 1947 by Bridgman, found that  $\Delta S_m$  and  $\Delta V_m$  always decrease with rising  $P_m$  and, indeed, extrapolate to zero at some finite high pressure, a criterion of a critical point. The behavior of He then appears to be anomalous, at least up to 3555 kg cm<sup>2</sup>. The continued rise with pressure of  $\Delta S_m$  is incompatible with the possibility of a critical point between solid and fluid. Since the question of a critical point in melting curves has yet to be resolved, it is interesting to extrapolate the He melting data to higher pressures than were measured.

An expression for  $\Delta S_m$  at high pressures can be derived in terms of  $P_m$  by combining Eqs. (1) and (3). When  $d\Delta S_m/dP_m$  is set equal to zero, one finds the solutions  $P_m = 4219$  kg/cm<sup>2</sup> and  $P_m = 3628$  kg/cm<sup>2</sup> for He<sup>3</sup> and He<sup>4</sup>, respectively.